

1D Coaxial Photonic Crystal - Superluminal



J. N. Munday and W. M. Robertson have reported observing: *Negative group velocity pulse tunneling through a coaxial photonic crystal*, Applied Physics Letters, Sept. 2002. In the following, we analyze this situation using Mathcad.

When a signal traverses an impedance boundary, it experiences a phase shift and a partial reflection that can be calculated from the generalized optical Fresnel coefficients of reflection r ($|V_{\text{reflected}}/V_{\text{incident}}|$) and transmission t ($|V_{\text{transmitted}}/V_{\text{incident}}|$). This can be generalized to coaxial cables using characteristic impedance, where z_i , z_r and z_t are the impedance of the incident, refelected and transmitted media.

$$r = \frac{z_i - z_t}{z_i + z_t} \quad t = \frac{2 \cdot z_i}{z_i + z_t}$$

Periodic variation in the impedance of a medium can produce destructive interference for some wavelengths. The phase accumulated throughout the crystal changes rapidly with frequency, especially near the band gap.

A unit cell consists of two coax segments, one of 50 ohm RG-58/U and one with 75 ohm RG-59/U. Each segment has the same phase velocity $0.66c$ and length 8 ft. As a result of impedance mismatch, 20% of the field is reflected at each interface. 12 unit cells of total length 120 m are used. A deep stop gap between 18 and 23 MHz occurs. Outside gap the attenuation is 25 - 35 dB/km. Stop occurs when path length of unit cell is multiple of $1/2$.

To calculate dispersive properties and group velocity, the effective index theory is used. The theory says phase shift and scattering loss of the electric field is from an effective complex index of refraction. the real part of the index, n_r , is obtained from the overall phase shift ϕ accumulated throught the crystal of length D . t is the complex coefficient of electric field transmission over the whole crystal and $m = 0$

$$n := 0 \quad \phi = \arctan\left(\frac{\text{Im}(t)}{\text{Re}(t)}\right) + n \cdot \pi \quad n_r(\omega) = \frac{c \cdot \phi}{\omega \cdot D} \quad f_{\text{stop}} = \frac{v}{2d}$$

We transmit a sinusoidal carrier with a Gaussian shaped pulse envelope. For carrier of 5 - 15 MHz pulse duration was scaled from 6 to 2 us, keeping the number of cycles within the envelope constant at 30 while varying the bandwidth fro 0.15 to 0.45 MHz.

COAXIAL CABLE MODEL

Physical Constants

Number of Sections, N 158 := 20·m 159 := 15·m

$$c := 299792458 \cdot \frac{\text{m}}{\text{sec}} \quad \mu_0 := 4 \cdot \pi \cdot 10^{-7} \cdot \frac{\text{newton}}{\text{amp}^2} \quad \epsilon_0 := 8.854187817 \cdot 10^{-12} \cdot \frac{\text{farad}}{\text{m}}$$

$$\rho_{\text{Cu}} := 1.673 \cdot 10^{-6} \cdot \Omega \cdot \text{cm} \quad \sigma := \rho_{\text{Cu}}^{-1} \quad \mu_r := 1 \quad \mu := \mu_0 \cdot \mu_r \quad \lambda_0(f) := \frac{c}{f}$$

RG-58/U (50 Ω) and RG-59/U (75 Ω) Coaxial Cable Data:

ϵ_r is the dielectric constant Ydb is attenuation in dB

$$\text{Crg58U}_{\text{len}} := 95 \cdot \frac{\text{pf}}{\text{m}} \quad a := \frac{1}{2} \cdot 0.9 \text{mm} \quad b := \frac{1}{2} \cdot 2.9 \text{mm} \quad \epsilon_r := 2.29 \quad v_{\text{c}} := 0.66 \quad 158 := 8 \cdot \text{ft}$$

$$\text{Crg59U}_{\text{len}} := 68.6 \cdot \frac{\text{pf}}{\text{m}} \quad a_{59} := \frac{1}{2} \cdot 0.81 \text{mm} \quad b_{59} := \frac{1}{2} \cdot 3.66 \text{mm} \quad \text{dB}_{100\text{m}} \quad 5,55,500\text{MHz} \quad 0.86,2, 5.7$$

$$\text{Crg62U}_{\text{len}} := \frac{\text{pf}}{\text{m}} \quad a_{59} := \frac{1}{2} \cdot 0.81 \text{mm} \quad b_{59} := \frac{1}{2} \cdot 3.66 \text{mm} \quad v_{\text{c}59} := 0.66 \quad 159 := 8 \cdot \text{ft}$$

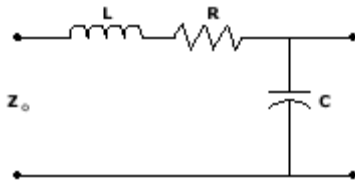
$$\text{Ydb}_{\text{Rg58}}(f_{\text{MHz}}) := 1.292 + 1.537 \cdot \sqrt{f_{\text{MHz}}} + 0.0157 \cdot f_{\text{MHz}} \quad \tan \delta := 0.0004 \quad v_{\text{c}62} := 0.85 \quad 162 := 7.97 \cdot \text{m}$$

$$\epsilon_{58} := \epsilon_0 \cdot \epsilon_r \quad d := 158 + 159$$

$$\epsilon_{59r} := \frac{1}{v_{\text{c}59}^2} \quad f_{\text{gap}} := \frac{v_{\text{c}} \cdot c}{2 \cdot d} \quad \epsilon_{59} := \epsilon_0 \cdot \epsilon_{59r}$$

Coaxial Cable Model

Conductor Radii a and b



$$C58_{Len} := \frac{2 \cdot \pi \cdot \epsilon_{58}}{\ln\left(\frac{b}{a}\right)}$$

$$C59_{Len} := \frac{2 \cdot \pi \cdot \epsilon_{59}}{\ln\left(\frac{b_{59}}{a_{59}}\right)}$$

$$\delta_{skin}(\omega) := \frac{1}{\sqrt{\pi \cdot \frac{\omega}{2 \cdot \pi} \cdot \mu_0 \cdot \mu_r \cdot \sigma}}$$

$$L58_{Len} := \frac{\mu}{2 \cdot \pi} \cdot \ln\left(\frac{b}{a}\right)$$

$$L59_{Len} := \frac{\mu}{2 \cdot \pi} \cdot \ln\left(\frac{b_{59}}{a_{59}}\right) \quad \sqrt{\frac{1}{\epsilon_r}} = 0.661$$

$$R58_{Len}(\omega) := \frac{a + b}{2 \cdot \pi \cdot a \cdot b \cdot \delta_{skin}(\omega) \cdot \sigma}$$

$$R59_{Len}(\omega) := \frac{a_{59} + b_{59}}{2 \cdot \pi \cdot a_{59} \cdot b_{59} \cdot \delta_{skin}(\omega) \cdot \sigma} \quad \epsilon_r \lambda = \left(\frac{\lambda_{coax}}{\lambda_0}\right)^2$$

p= r Return Loss = $10 \log_{10} \left(\frac{1}{1-p^2} \right)$ dB

$$\lambda(f) := c \cdot (\sqrt{\epsilon} \cdot f)$$

$$v = \frac{1}{\sqrt{L_{len} \cdot C_{len}}} = \frac{2 \cdot \text{Length}}{\Delta t}$$

$$f_{gap} = 2.029 \times 10^7 \frac{1}{s}$$

Characteristic Impedance

(a is conductor and b is dielectric radii, Llen and Clen are L and C per length.):

$$Z_o(\omega) := \sqrt{\frac{R + j \cdot \omega \cdot L}{G + j \cdot \omega \cdot C}}$$

$$Z58_{olc} := \sqrt{\frac{L58_{Len}}{C58_{Len}}}$$

$$Z_o := 138 \cdot \sqrt{\frac{1}{\epsilon_{59r}}} \cdot \log\left(\frac{b_{59}}{a_{59}}\right)$$

$$Z_o = 59.657$$

Attenuation, Phase and Propagation Constants: α , β , γ

$$Z58_{olc} = 46.36 \Omega$$

$$\gamma(\omega) = \alpha(\omega) + j \cdot \beta(\omega) \quad \gamma(\omega) := \sqrt{(R + j \cdot \omega \cdot L)(G + j \cdot \omega \cdot C)} \quad \alpha(\omega) := \text{Re}(\gamma(\omega)) \quad \beta(\omega) := \text{Im}(\gamma(\omega))$$

$$E(\omega, x) = E_0 \cdot e^{-(\alpha(\omega) + j \cdot \beta(\omega)) \cdot x} \quad V(z) = V_{incident} \cdot e^{-\gamma \cdot z} + V_{reflected} \cdot e^{\gamma \cdot z} \quad \alpha_s(f) := 2 \cdot \beta(2 \cdot \pi \cdot f)$$

$$\text{AbsorbCoef}(f) := 10^{-7} \cdot \left(2 \cdot \pi \cdot \frac{\sqrt{f}}{m} \right) \quad \kappa(f) := \sqrt{\mu_0 \cdot \epsilon_0 \cdot \epsilon_r} \cdot 2 \cdot \pi \cdot \frac{f}{\text{sec}} + j \cdot \frac{\text{AbsorbCoef}(f)}{2} \text{ wave number}$$

$$E(x) = E_0 \cdot e^{j \cdot (\kappa \cdot x + \phi)}$$

Material Loss tangent, $\tan \delta$ equals $jG_{\omega C}$ Thus $G = j \cdot \omega \cdot C \cdot \tan \delta = 2 \cdot \frac{\pi \cdot \epsilon_0 \cdot \epsilon_r \cdot \text{length} \cdot \omega \cdot \tan \delta}{\ln\left(\frac{b}{a}\right)}$

$$G58(\omega) := j \cdot \omega \cdot Crg58U_{len} \cdot \tan \delta$$

$$G59(\omega) := j \cdot \omega \cdot Crg59U_{len} \cdot \tan \delta$$

$$Z58_o(f) := \sqrt{\frac{R58_{Len} \left(\frac{2 \cdot \pi \cdot f}{\text{sec}} \right) + j \cdot \frac{2 \cdot \pi \cdot f}{\text{sec}} \cdot L58_{Len}}{G58 \left(\frac{2 \cdot \pi \cdot f}{\text{sec}} \right) + j \cdot \frac{2 \cdot \pi \cdot f}{\text{sec}} \cdot Crg58U_{len}}} \quad Z59_o(f) := \sqrt{\frac{R59_{Len} \left(\frac{2 \cdot \pi \cdot f}{\text{sec}} \right) + j \cdot \frac{2 \cdot \pi \cdot f}{\text{sec}} \cdot L59_{Len}}{G59 \left(\frac{2 \cdot \pi \cdot f}{\text{sec}} \right) + j \cdot \frac{2 \cdot \pi \cdot f}{\text{sec}} \cdot Crg59U_{len}}}$$

Reflection and Transmission Coefficients, t and r

From Maxwells Equations: Plane incident E-M wave traveling from medium A to B

$$t_{AB}(f) := \frac{2 \cdot Z59_o(f)}{Z58_o(f) + Z59_o(f)}$$

$$t_{BA}(f) := \frac{2 \cdot Z58_o(f)}{Z58_o(f) + Z59_o(f)}$$

$$r_{AB}(f) := \frac{Z59_o(f) - Z58_o(f)}{Z58_o(f) + Z59_o(f)}$$

$$r_{BA}(f) := \frac{Z58_o(f) - Z59_o(f)}{Z58_o(f) + Z59_o(f)}$$

Structure of Lengths of N unit cells starting with Len58 N := 12

$$j := 0..N \quad \text{Len}_{58_j} := 158 \quad \text{Len}_{58_0} := 158 \quad \text{Len}_{59_j} := 159 \quad \text{Len}_{59_0} := 0 \cdot m \quad D := N(d)$$

$$M_{BA}(f, j) := \begin{pmatrix} \frac{e^{-j \cdot \kappa(f) \cdot \text{Len}_{58_j}}}{t_{BA}(f)} & \frac{r_{BA}(f) \cdot e^{j \cdot \kappa(f) \cdot \text{Len}_{58_j}}}{t_{BA}(f)} \\ \frac{r_{BA}(f) \cdot e^{-j \cdot \kappa(f) \cdot \text{Len}_{58_j}}}{t_{BA}(f)} & \frac{e^{j \cdot \kappa(f) \cdot \text{Len}_{58_j}}}{t_{BA}(f)} \end{pmatrix}$$

$$M_{AB}(f, j) := \begin{pmatrix} \frac{e^{-j \cdot \kappa(f) \cdot \text{Len}_{59_j}}}{t_{AB}(f)} & \frac{r_{AB}(f) \cdot e^{j \cdot \kappa(f) \cdot \text{Len}_{59_j}}}{t_{AB}(f)} \\ \frac{r_{AB}(f) \cdot e^{-j \cdot \kappa(f) \cdot \text{Len}_{59_j}}}{t_{AB}(f)} & \frac{e^{j \cdot \kappa(f) \cdot \text{Len}_{59_j}}}{t_{AB}(f)} \end{pmatrix}$$

Field Amplitudes, E, for Transmitted and Reflected Waves

$$E(f) := \left[\prod_{j=1}^N \left(M_{AB}(f \cdot 10^6, j) \cdot M_{BA}(f \cdot 10^6, j-1) \right) \right] \cdot \begin{pmatrix} 1 \\ 0 \end{pmatrix} \quad \text{Phase } \phi - \text{Effective Index Theory}$$

$$\phi(f_{\text{MHz}}) := -\arg \left[(E(f_{\text{MHz}})_0)^{-1} \right]$$

$$T_s(f_{\text{MHz}}) := \left| (E(f_{\text{MHz}})_0)^{-1} \right|$$

$$\text{Phz} := \begin{cases} G \leftarrow 0 \\ P_1 \leftarrow \phi \left(\frac{1}{10} \right) \\ \text{Pold} \leftarrow P_1 \\ \text{for } n \in 2..150 \\ \left| \begin{array}{l} P_n \leftarrow \phi \left(\frac{n}{10} \right) \\ G \leftarrow G + 1 \text{ if } P_n > \text{Pold} + 2 \\ \text{Pold} \leftarrow P_n \\ P_n \leftarrow P_n - G \cdot \pi \end{array} \right. \\ \text{P} \end{cases}$$

Group time lag, tg, Group Velocity, vg

$$\tau_g(f) := - \left(\frac{1}{2 \cdot \pi \cdot \frac{10^6}{\text{sec}}} \cdot \frac{\phi(f + 0.01) - \phi(f - 0.01)}{0.02} \right)$$

$$v_g(f) := \frac{D}{c \cdot \tau_g(f)} \quad \text{WRITEPRN("PhaseBrk")} := \text{Phz}$$

$$\phi_{\text{break}} := \text{READPRN("PhaseBrk.prn")}$$

Calculate Complex Vector Arrays to Store Results for Plotting

$$n := 1..250 \quad \text{freq}_n := \frac{n}{10} \quad E_{f_n} := E(\text{freq}_n) \quad T_{sf_n} := \left| \left[(E_{f_n})_0 \right]^{-1} \right| \quad \phi_{f_n} := -\arg \left[\left[(E_{f_n})_0 \right]^{-1} \right]$$

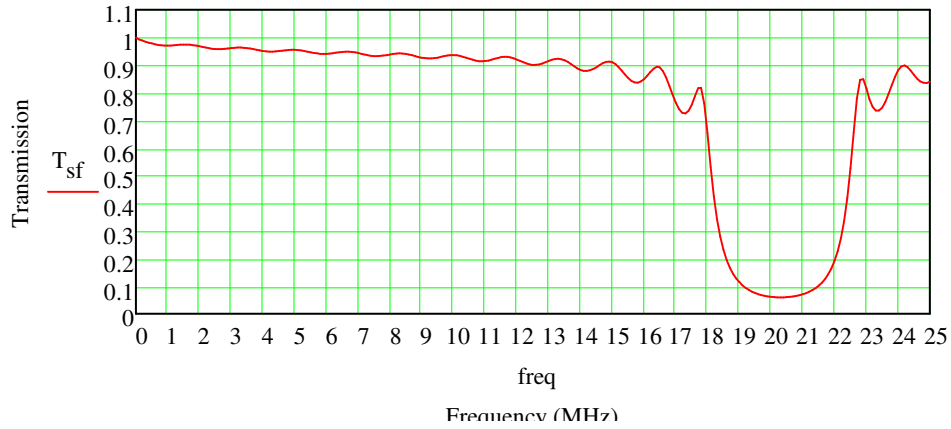
$$\tau_{gf_n} := - \left(\frac{1}{2 \cdot \pi \cdot \frac{10^6}{\text{sec}}} \cdot \frac{\phi_{f_n} - \phi_{f_{n-1}}}{0.1} \right) \quad v_{gf_n} := \frac{D}{c \cdot \tau_{gf_n}}$$

$$c_{\text{vacuum}} := 1$$

$$f_{\text{gap}} = 2.029 \times 10^7 \frac{1}{s}$$

$$T_{sf_0} := 1 \quad \text{dbT} := \overrightarrow{(10 \cdot \log(T_{sf}))}$$

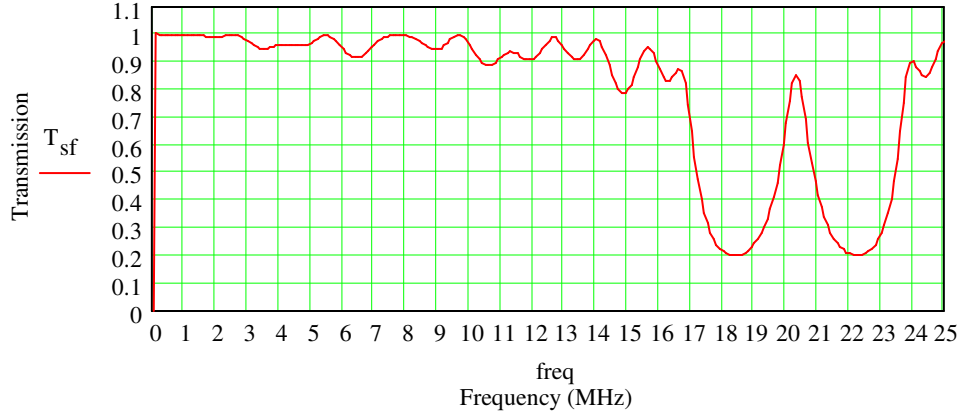
Photonic Gap Transmission Spectra



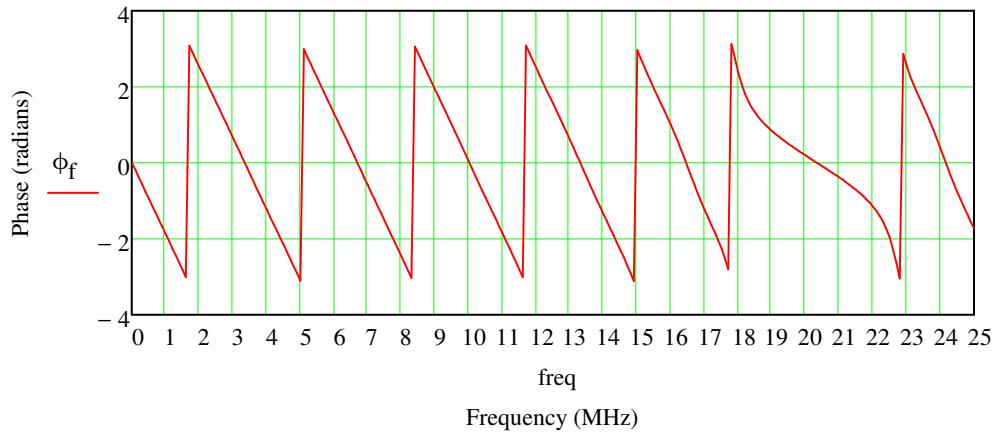
$$\text{match}(\min(T_{sf}), T_{sf})^T \cdot 0.1 = (20.1 \ 20.2 \ 20.3 \ 20.4 \ 20.5)$$

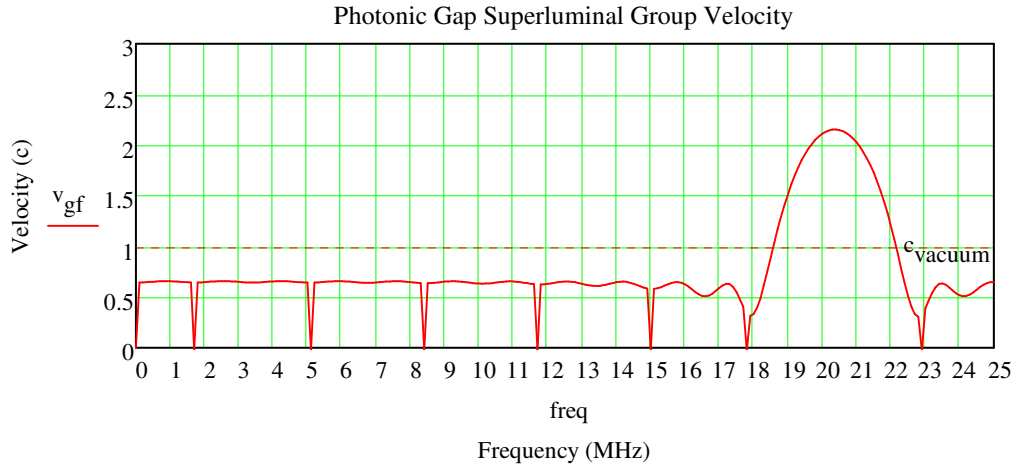
Result (picture) for $\text{Len}_{58_2} := 4.158$ $\text{Len}_{58_\epsilon} := 4.158$

Photonic Gap Transmission Spectra

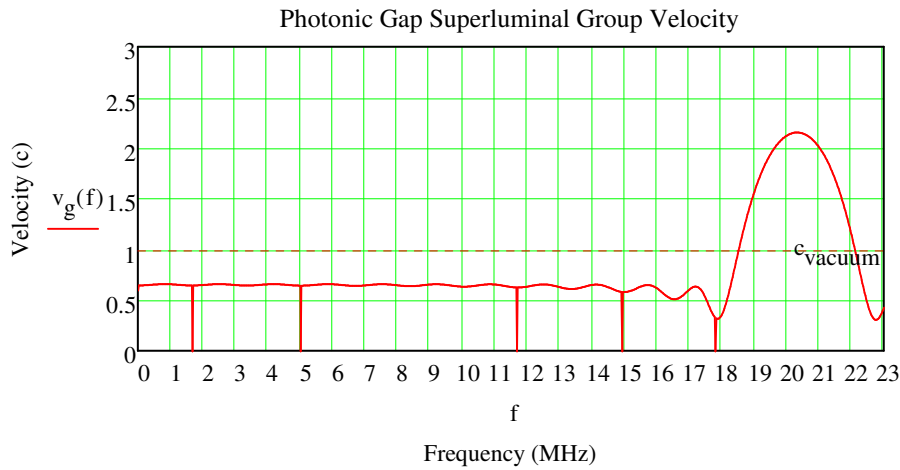


Photonic Gap Phase Lag



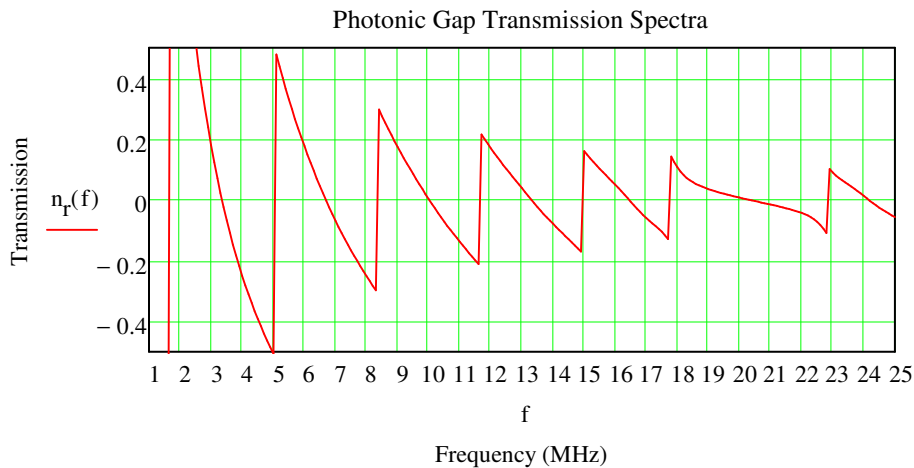


Below are exact solutions. Evaluation was disabled because they take a long time to calculate each function and then repaint each of the graphs



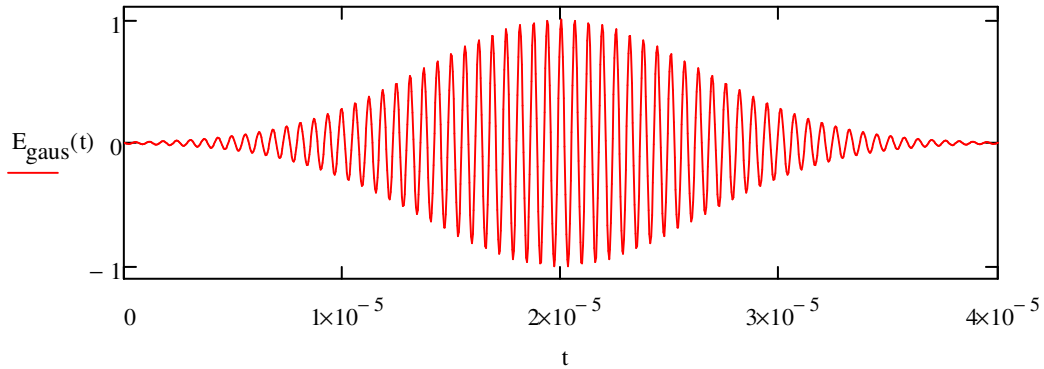
$$\Delta t := \frac{12 \cdot (158 + 159)}{0.66 \cdot c} \quad \Delta t = 295.768 \text{ nsec}$$

$$n_r(f) := \frac{c \cdot \phi(f)}{2\pi \cdot \frac{f}{\text{sec}} \cdot 10^6 \cdot D} \quad n_r(1) = -1.526 \quad f := 1, 1.1.. 25$$



Type of Pulse Wave Packet for Transmission into Photonic Crystal Cable

$$\omega_0 := 1 \cdot 10^7 \quad \tau := 10 \cdot \frac{2 \cdot \pi}{\omega_0} \quad \phi_g := 2 \cdot 10^{-5} \quad E_{\text{gaus}}(t) := \text{Re} \left[e^{\left[\frac{-(t-\phi_g)^2}{2 \cdot \tau^2} - j \cdot \omega_0 \cdot (t-\phi_g) \right]} \right]$$



References

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